



*LSU Relativity Seminar
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Constant-expansion surfaces for finite-distance angular momentum estimates in numerical relativity

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OUTLINE

- **MOTIVATION:** estimating angular momentum in simulations of strongly gravitating systems;
 - For establishing balance laws, study transfer of angular momentum in non-vacuum scenarios;
 - For error control in numerical settings [Baker2007, Pollney2007];
 - For initial data [Caudill2006];
- **BACKGROUND:**
 - Angular momentum ambiguities in General Relativity;
 - Measuring angular momentum in Numerical Relativity;
 - Constant-Expansion (CE) surfaces;
- **RESULTS**
- **CONCLUSIONS**

Conservation laws for matter

- Noether procedure for the construction of conserved currents [Szabados2009 and references];
- In its possibly simplest form: test mass in free fall:

$$\frac{du^a}{d\tau} + \Gamma_{bc}^a u^b u^c = 0$$

$$\xi_a \left(\frac{du^a}{d\tau} + \Gamma_{bc}^a u^b u^c \right) = \frac{d}{d\tau} (\xi_a u^a) - u^a u^b \xi_{a;b}$$

- There is a conserved quantity for each Killing vector field:
 - In flat spacetimes, Poincaré group:
 - ▶ Translations → energy-momentum
 - ▶ Rotations → angular momentum
 - ▶ Lorentz boosts → center-of-mass conservation
 - In general spacetimes? [Harte2008]

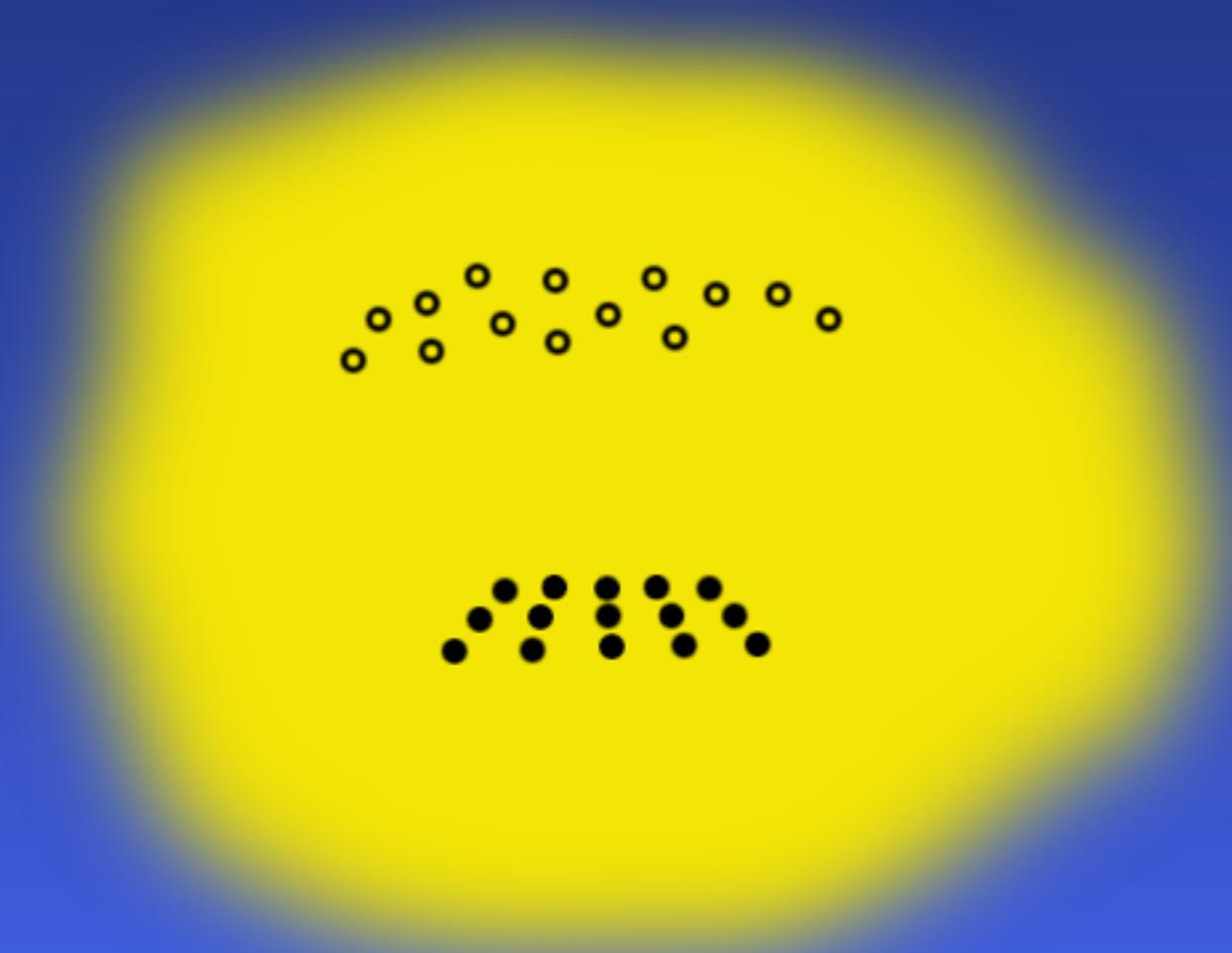
Conservation laws for the gravitational field

- General Relativity: apply to the gravitational field?

PROBLEMS	SOLUTIONS
<p>LOCALIZATION: “local” energy-momentum and angular-momentum of the gravitational field are not observable (local flatness).</p>	Global approach (spatial and null infinity)
	Quasi-local approach (extended but finite regions)
<p>SYMMETRY & CONSERVATION: in a generic gravitational field, which symmetries? How to define translations? Rotations?</p>	Asymptotic symmetries
	Quasi-local symmetries
	Approximate symmetries
<p>ORIGIN: Choose an appropriate origin with respect to which angular momentum is calculated (center of mass problem in GR). Important if one wants to compare different spacetimes!</p>	Good cuts, nice sections, CE surfaces?

Additional complications

- In numerical simulations, however, it is not trivial to compute these quantities:
 - Only **finite-distance surfaces** are usually available - how can one calculate the limit to spatial or null infinity?
 - **Gauge unknown**, potentially non-trivial, not modifiable - how can one ensure that these corresponds to the correct chart/tetrad choice required by the global formalisms?



MEASURING ANGULAR MOMENTUM IN NUMERICAL RELATIVITY

Solutions

APPROACH	SOLUTION
<p>Common approach in numerical simulations</p>	<p>Use ADM angular momentum</p> $J_{\text{ADM}} = \frac{1}{8} \oint \varphi^a R^b K_{ab} d^2V$ <p>or Bondi angular momentum</p> $J_{\text{Bondi}}^i = \frac{1}{16\pi} \text{Re} \left\{ \oint [\bar{m}^i (2\Psi_1 - 2\sigma\bar{\sigma} - \bar{\sigma}(\sigma\bar{\sigma})) + \ell^i (\Psi_2 + \sigma\bar{\sigma} - \bar{\sigma}^2\sigma)] d^2V \right\}$ <p>with chart and tetrad build from simulation's coordinates; extract on a set of constant coordinate radius surfaces and extrapolate to infinity, but using:</p> $\Psi_4 = -\ddot{\bar{\sigma}}$
<p>Gallo2008, Deadman2008</p>	<p>Use the transformation law to a Bondi chart/tetrad to rewrite the Bondi expressions at finite distance and in generic coordinates.</p>
<p>Our test</p>	<p>Use the angular momentum from rotational Killing vector field</p> $J_{\text{Killing}} = \frac{1}{8} \oint \phi^a R^b K_{ab} d^2V$ <p>on CE surfaces, which select a radial coordinate, a tetrad and a center of mass.</p>

CONSTANT-EXPANSION SURFACES

Definitions

- In a 3+1 setting, given a 2-surface, a null tetrad can be constructed from the normal to the spatial hypersurfaces, the normal to the 2-surface on the spatial hypersurfaces and two vectors tangent to the 2-surface, using:

$$\begin{aligned}\ell^a &= \frac{T^a + R^a}{\sqrt{2}} \\ n^a &= \frac{T^a - R^a}{\sqrt{2}} \\ m^a &= \frac{\theta^a + i\varphi^a}{\sqrt{2}}\end{aligned}$$

- The surface's outgoing null expansion is defined as:

$$\Theta_{(\ell)} = q^{ab}\nabla_a\ell_b$$

$$q^{ab} = g^{ab} + \ell^a n^b + n^a \ell^b$$

- A 2-surface is said to be a CE surface if the expansion of its outgoing null normal is constant across it.
- On simple, exact spacetimes, the expansion is asymptotically monotonic with the coordinate radius, and can be used as a radial coordinate;

RESULTS

Setup

- Consider a Kerr spacetime in Kerr-Schild coordinates, and add a Lorentz boost and a traslation:

$$g_{ab} = \eta_{ab} + 2Hk_a k_b$$

$$H = \frac{Mr^3}{r^4 + a^2 z^2}$$

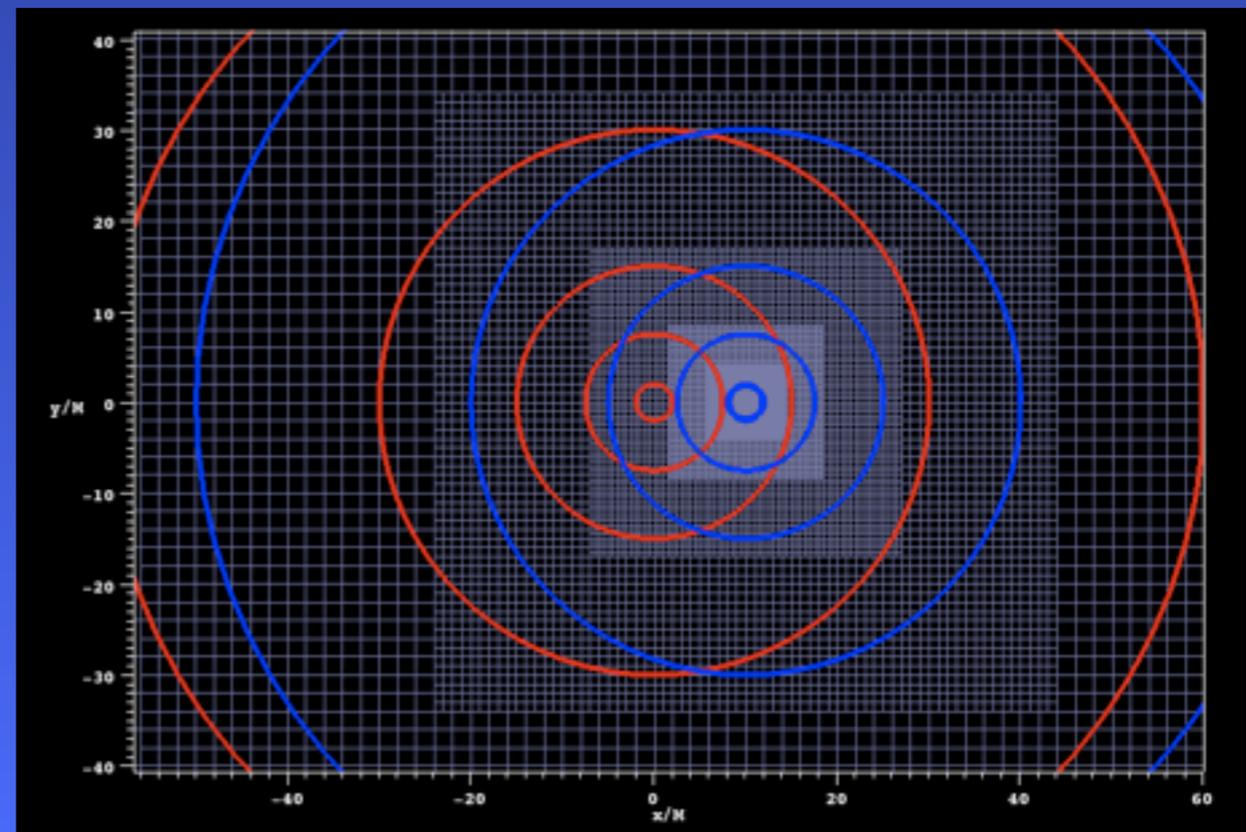
$$k_a dx^a = -\frac{r(xdx + ydy) - a(xdy - ydx)}{r^2 + a^2} - \frac{zdz}{r} - dt$$

$$t \rightarrow \gamma(t - \beta y)$$

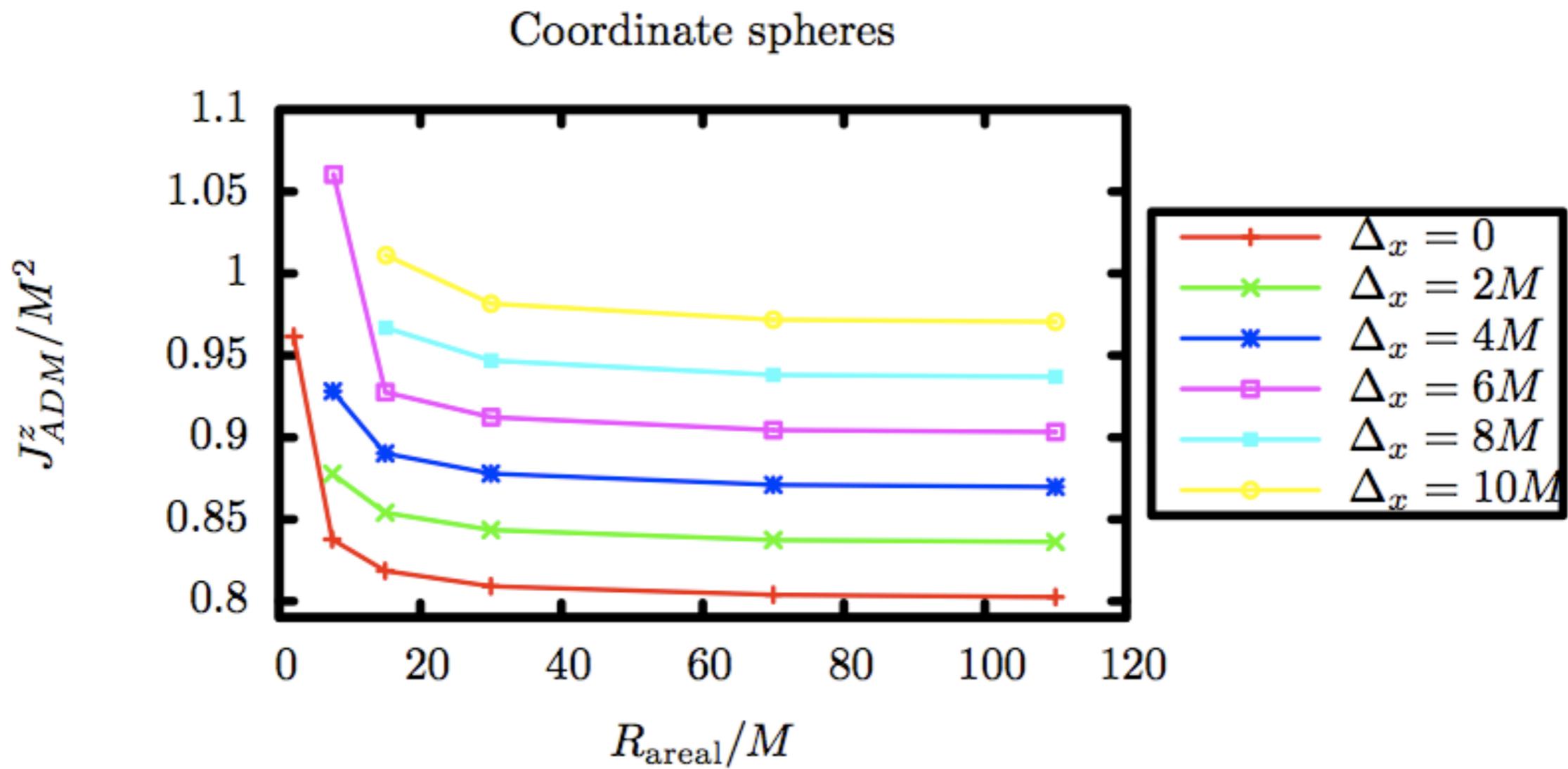
$$x \rightarrow x + \Delta$$

$$y \rightarrow \gamma(y - \beta t)$$

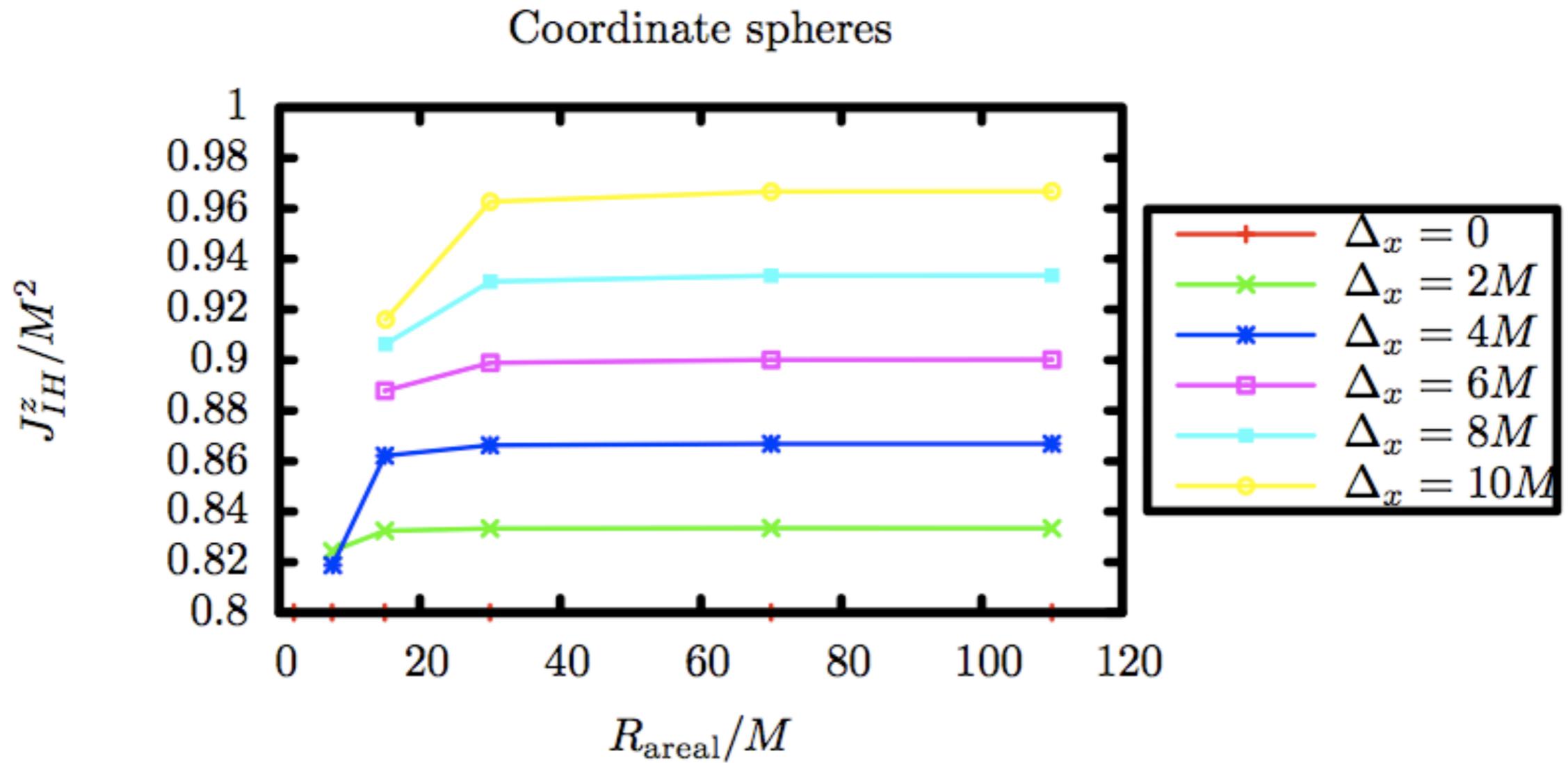
- Use a boost of 0.01668 and offset of (2, 4, 6, 8, 10)M.
- Construct two classes of surfaces: coordinate spheres and CE surfaces, located at areal radii of (2, 7.5, 15, 30, 60, 110)M;
- Integrate two types of angular momentum.



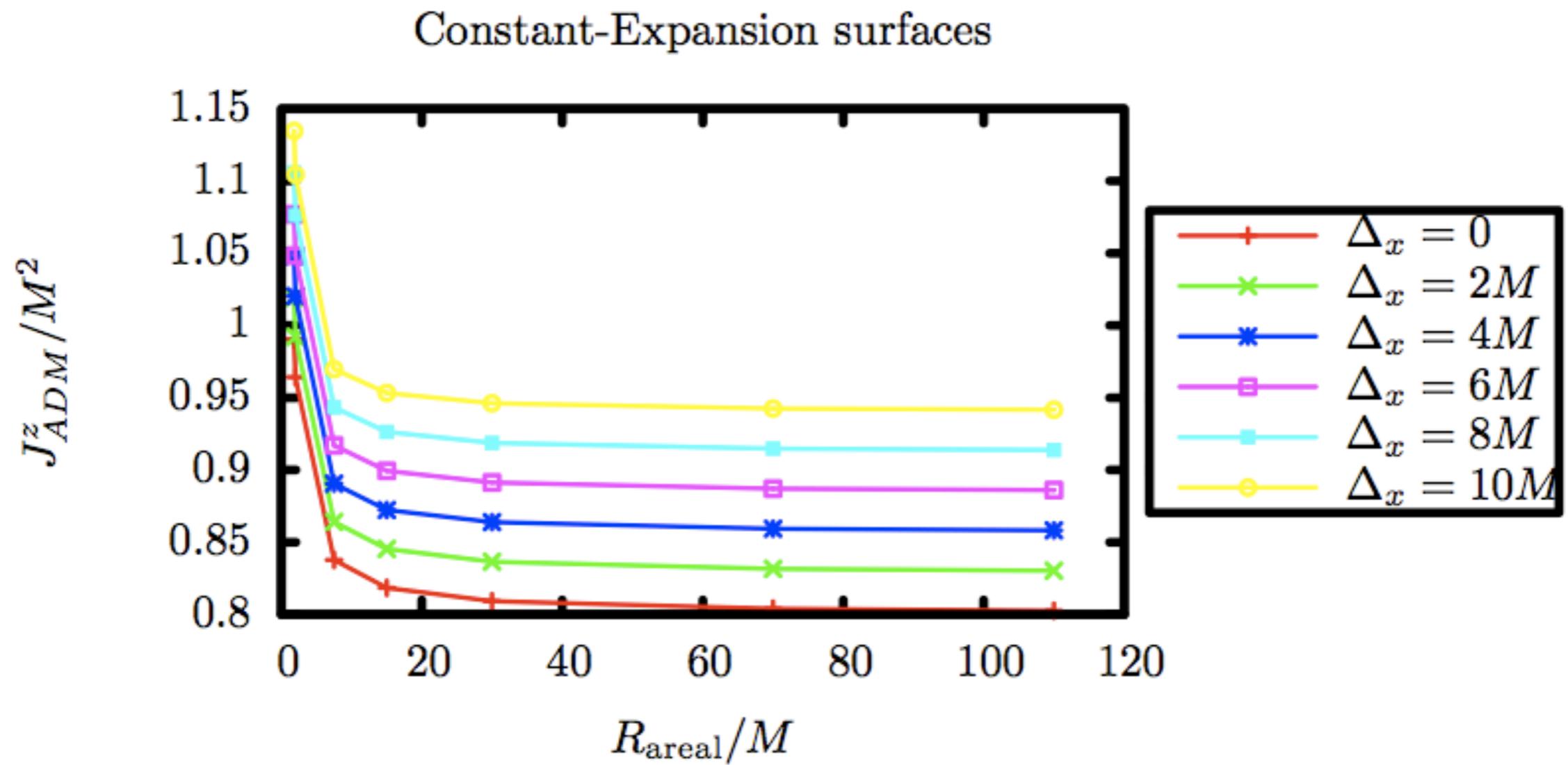
RESULTS



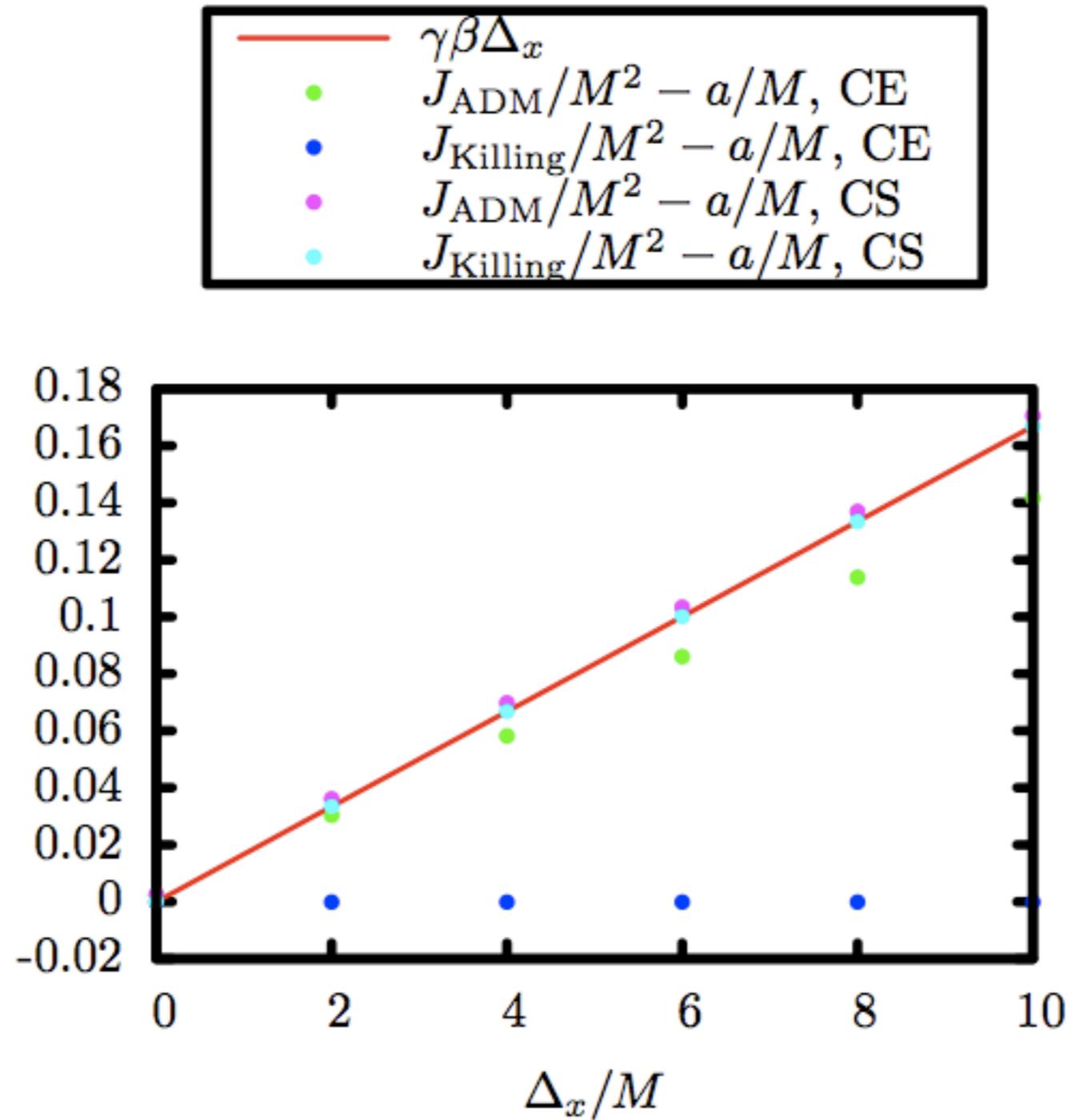
RESULTS



RESULTS



RESULTS



CONCLUSIONS

- **No natural way** to resolve the supertranslation ambiguity in a generic spacetime - no corresponding structure in the asymptotic isometry group, **ambiguity must be resolved by hand** with a sensible prescription;
- Equivalent to the problem of **finding a center of mass** in a generic, asymptotically flat spacetime;
- Angular momentum constructions at null infinity will always contain one such (more or less explicit) assumption: if the assumption about the center of mass is not the intended one, the angular momentum will have a **systematic error that does not vanish asymptotically**;
- For (slightly) boosted Kerr, the Killing construction calculates the angular momentum with respect to an observer corresponding to the integration surfaces; the ADM integrand yields unclear results.